

**MAGNETISM  
IN CARBON**

Amurri Alessandro

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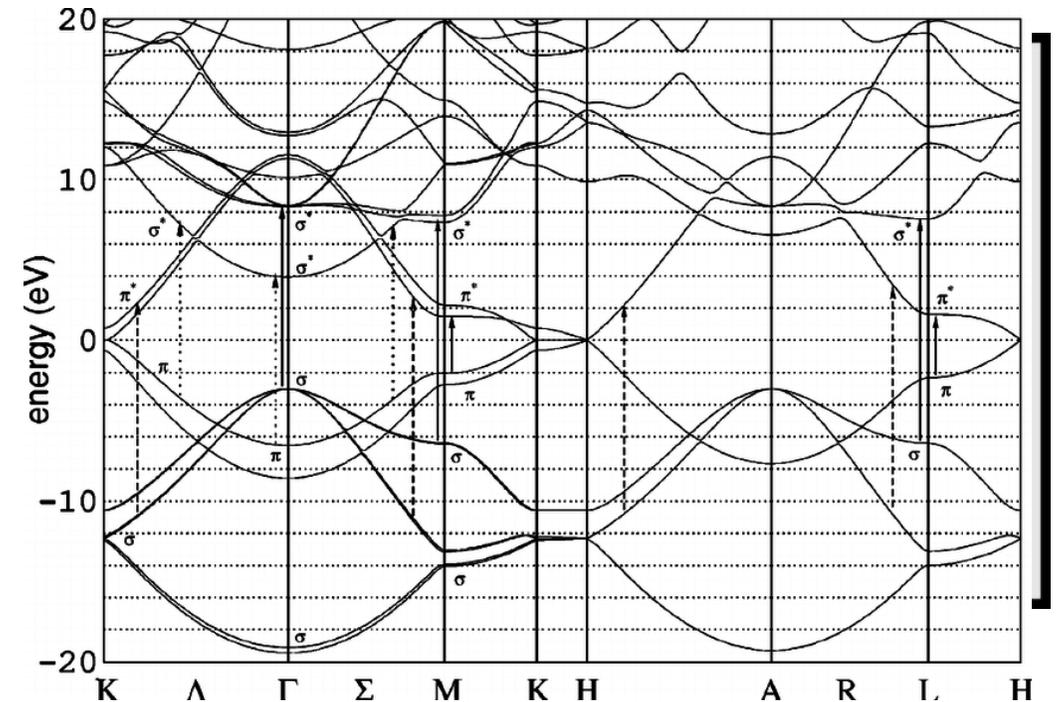
# INTRODUCTION

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# INGREDIENTS FOR MAGNETISM

- Conventional magnetism: d and f elements
- Partially filled and localized orbitals
- Strong exchange
- Broken spin degeneracy
- In graphite and graphene no unpaired electrons
- Strong overlap of orbitals
- No exchange
- Low DOS at Fermi level



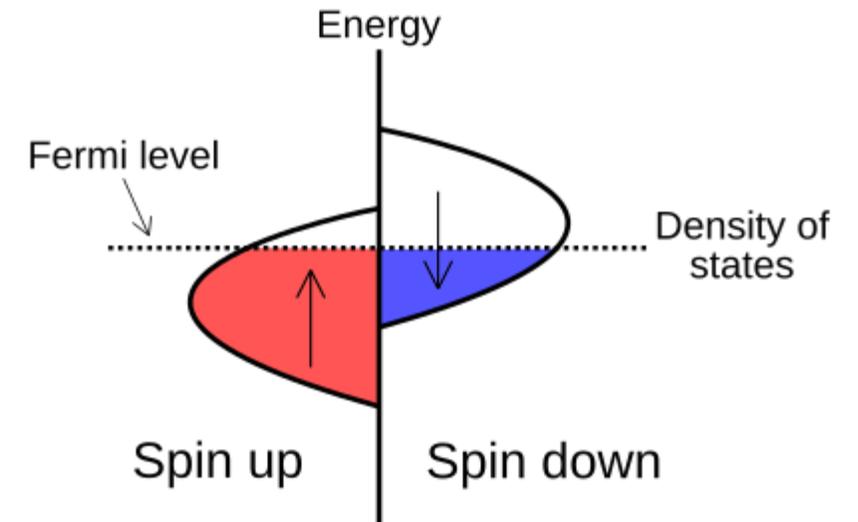
MARINOPOULOS, AG. et al. 2004 PHYSICAL REVIEW B

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# WHEN CAN P ELEMENTS BE MAGNETIC?

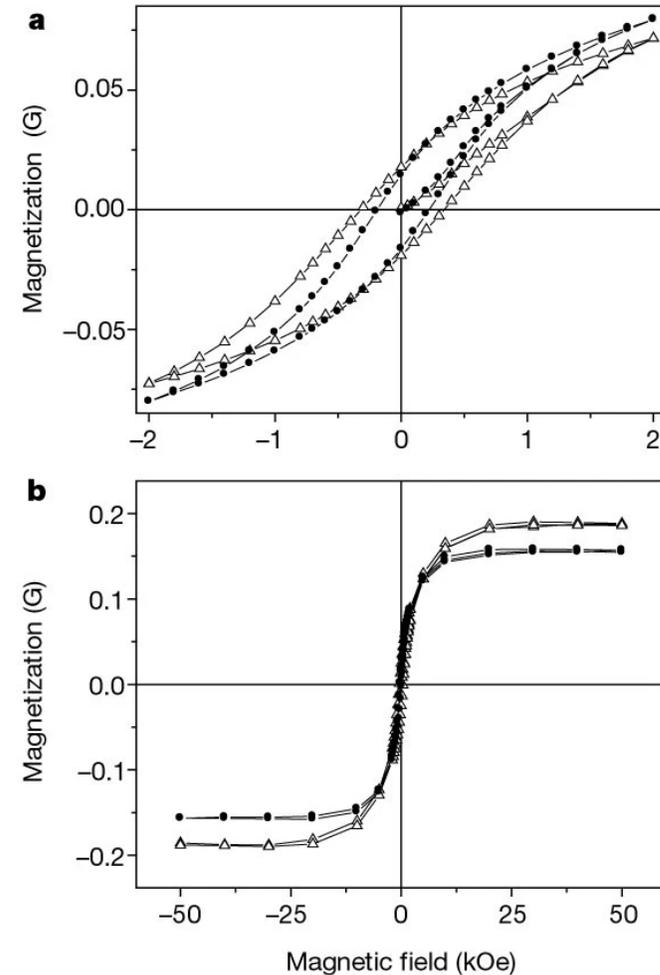
- Delocalized p electrons  $\rightarrow$  no magnetism
- Defects can induce a localization of  $\pi$  states
- Localization + electron-electron interaction  $\rightarrow$  unpaired spins
- Broken symmetry lifts spin degeneracy
- High DOS at Fermi level



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# FIRST (CONTROVERTIAL) EVIDENCE

- 2001 – Makarova et al., “Magnetic carbon” (Nature)
  - Claimed RT ferromagnetism in polymerized fullerenes
- The article was retracted
- Sparked interest in magnetism in carbon
- New evidences of defect-induced magnetism



T.L. Makarova et al. Nat. (2006) 413(6857)

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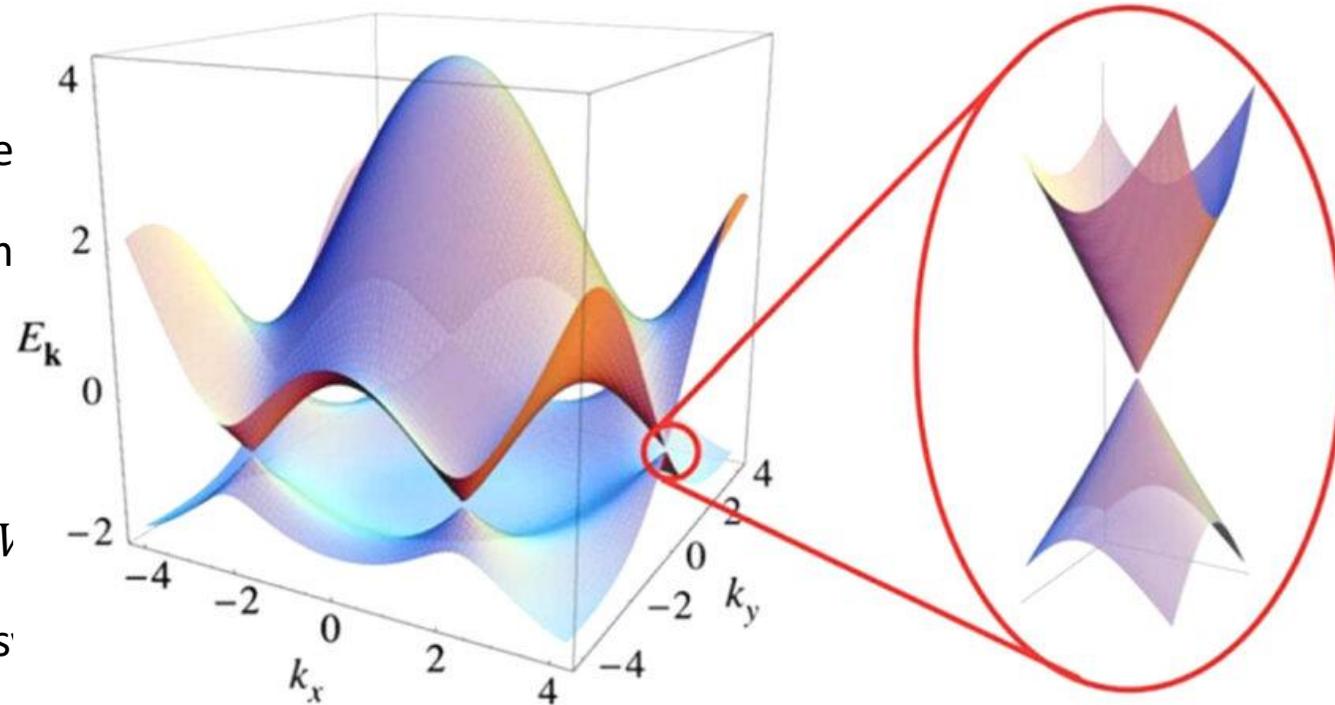
**THEORY**

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# THEORETICAL FRAMEWORK

- $\mathcal{H}_0$ : first ne
- Only  $\pi$ -sym
- $t \simeq -2,7eV$
- Half-filled s
- $\mathcal{H}_0$  exhibits electron-hole symmetry



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# THEORETICAL FRAMEWORK

$$\mathcal{H} = \mathcal{H}_0 + \mathcal{H}' = -t \sum_{\langle i,j \rangle, \sigma} [c_{i\sigma}^\dagger c_{j\sigma} + h.c.] + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

- $\mathcal{H}'$ : on-site Coulomb interactions between electrons
- $U > 0$ : Coulomb repulsion parameter of  $\pi$ -electrons
- Mean-field approximation is used

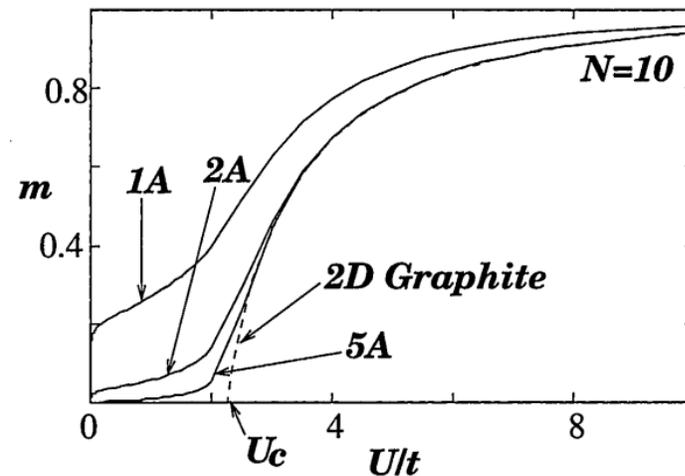
$$\mathcal{H}'_{mf} = U \sum_i [n_{i\uparrow} \langle n_{i\downarrow} \rangle + n_{i\downarrow} \langle n_{i\uparrow} \rangle - \langle n_{i\uparrow} \rangle \langle n_{i\downarrow} \rangle]$$

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# THEORETICAL FRAMEWORK

- Self-consistent solution  $M_i = \frac{\langle n_{i\uparrow} \rangle - \langle n_{i\downarrow} \rangle}{2} \rightarrow S = \sum_i M_i$
- The solution is controlled by the ratio  $U/t$



# THEORETICAL FRAMEWORK: COUNTING

## RULES

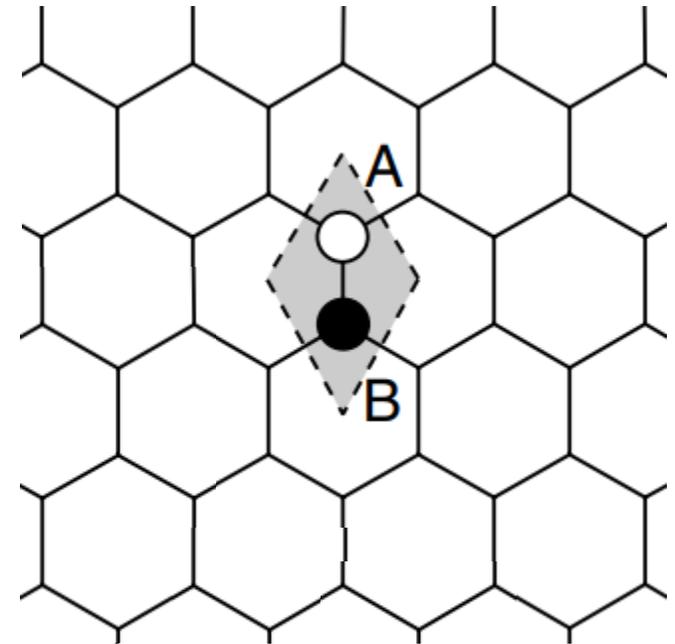
- From honeycomb lattice (benzenoid) group theory the number of zero energy states (ZES) is

$$\eta = 2\alpha - N$$

- Stoner criterion controls the onset of magnetism  $U\rho(E_f) > 1$

- Lieb's theorem determines the total spin

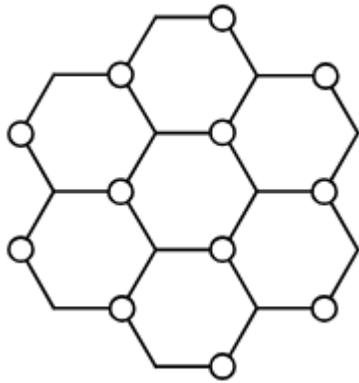
$$S = \frac{1}{2} |N_A - N_B|$$



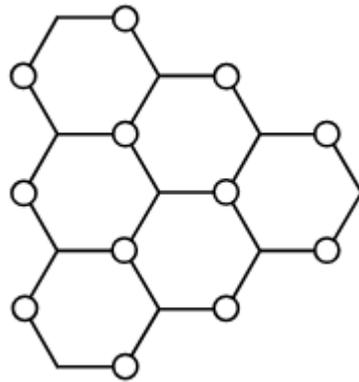
# THEORETICAL EXAMPLE: FINITE FRAGMENTS

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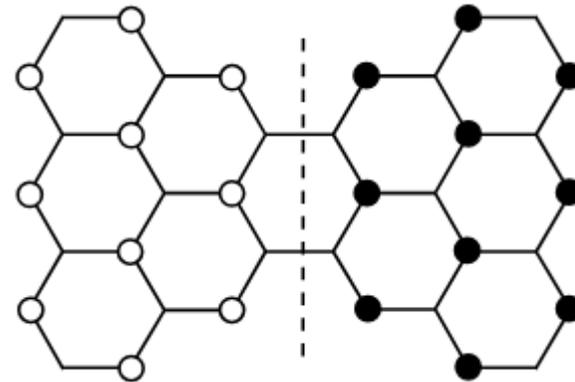
- Hexagonal finite graphene fragments



corenene



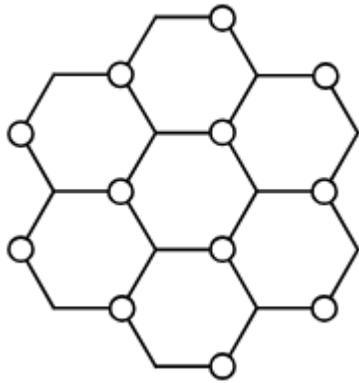
triangulane



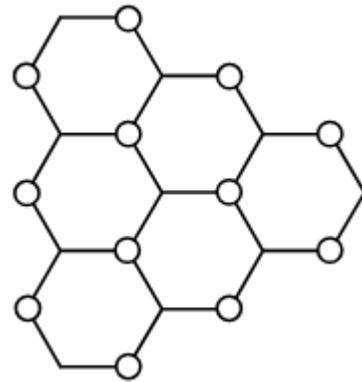
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# THEORETICAL EXAMPLE: FINITE FRAGMENTS

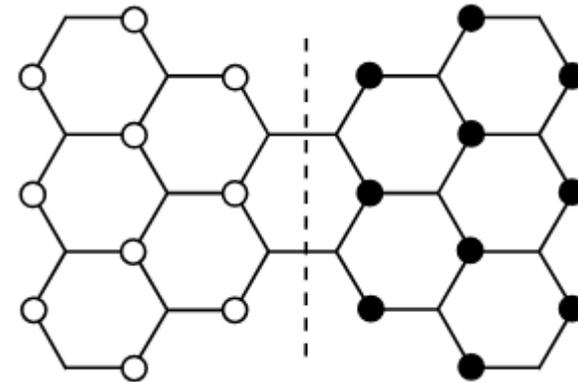
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$$\begin{aligned}N_A &= N_B = 12 \\ \alpha &= 12 \\ \eta &= 0 \\ S &= 0\end{aligned}$$



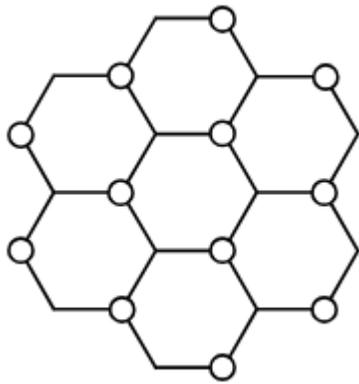
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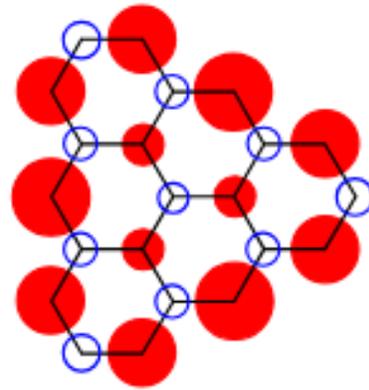
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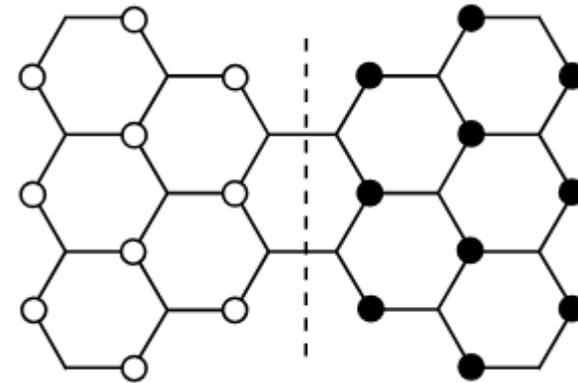
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corenene



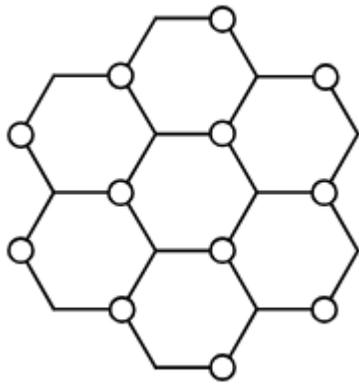
$$\begin{aligned} N_A - N_B &= 2 \\ \alpha &= 12 \\ \eta &= 2 \\ S &= 1 \\ m &= 2\mu_B \end{aligned}$$



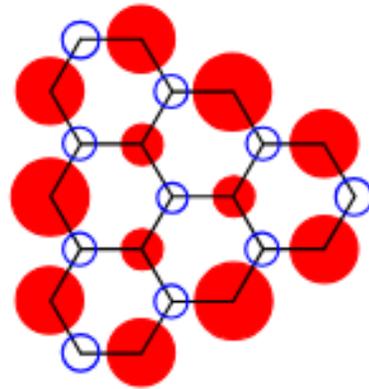
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# THEORETICAL EXAMPLE: FINITE FRAGMENTS

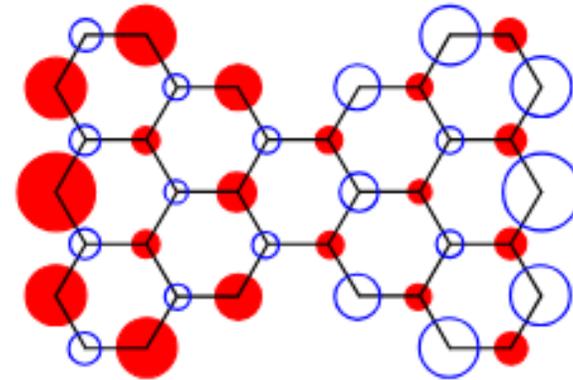
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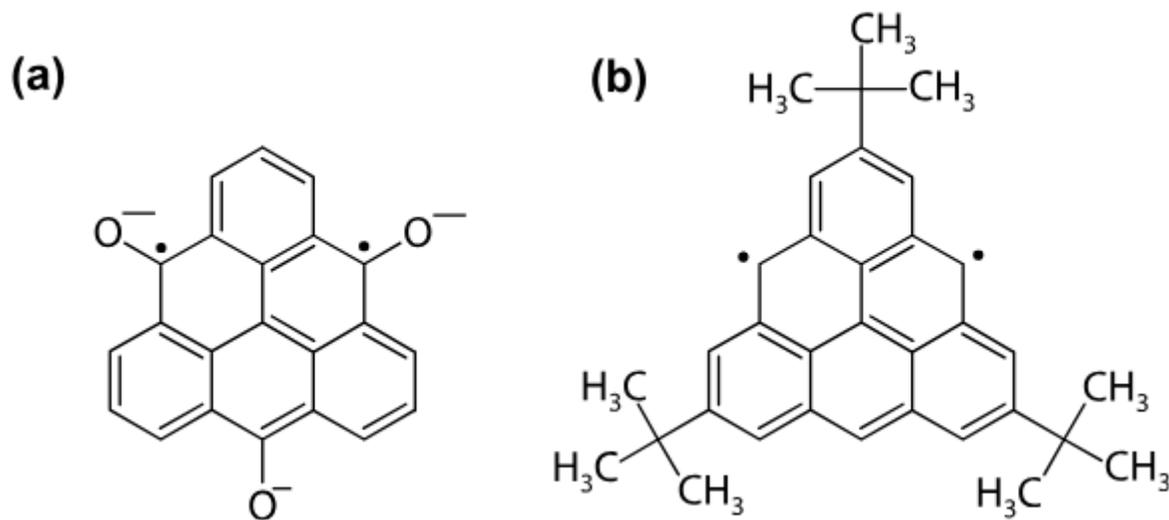
triangulane



$$\begin{aligned} N_A &= N_B = 19 \\ \alpha &= 20 \\ \eta &= 2 \\ S &= 0 \end{aligned}$$

# THEORETICAL EXAMPLE: FINITE FRAGMENTS

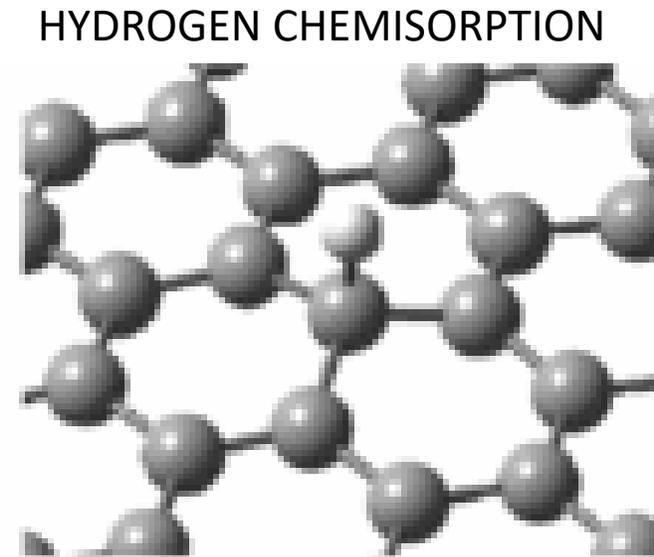
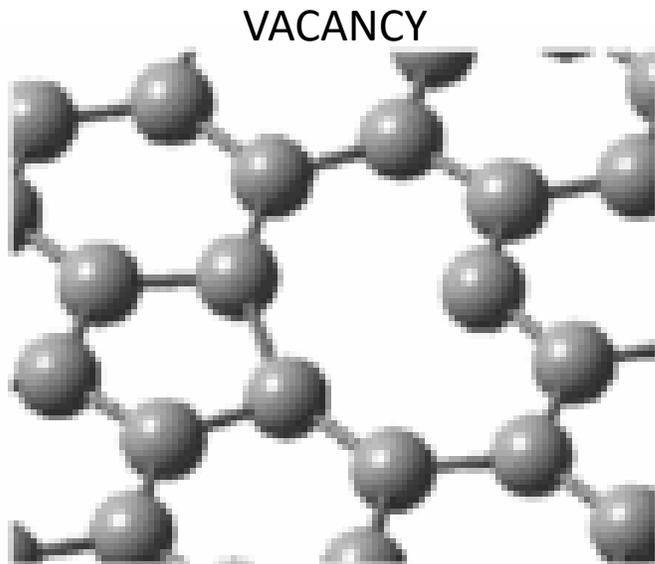
- Attempts to synthesize these molecule failed



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# DEFECTS INDUCED MAGNETISM

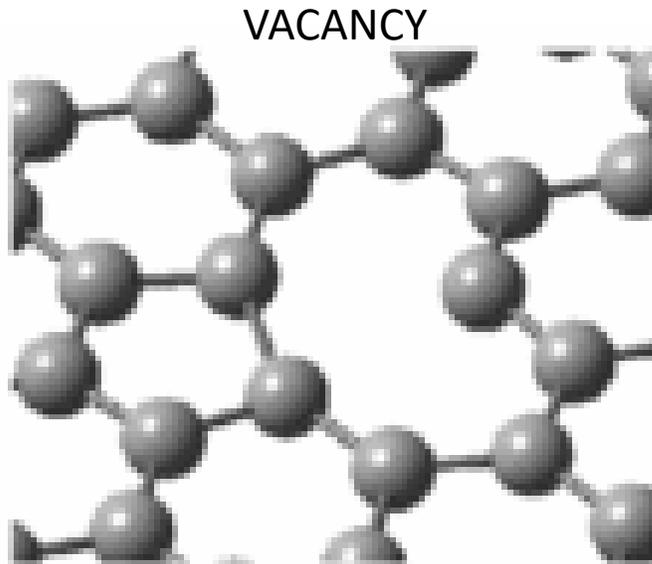
- Imbalance in sublattices leads to magnetism  $N_I \equiv |N_A - N_B| \geq 0$
- Simple defects in graphene



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# DEFECTS INDUCED MAGNETISM

- Imbalance in sublattices leads to magnetism  $N_I \equiv |N_A - N_B| \geq 0$
- Simple defects in graphene



- One carbon is removed
- A dangling bond is unsaturated
- Jahn-Teller distortion
- Imbalance in the sublattices

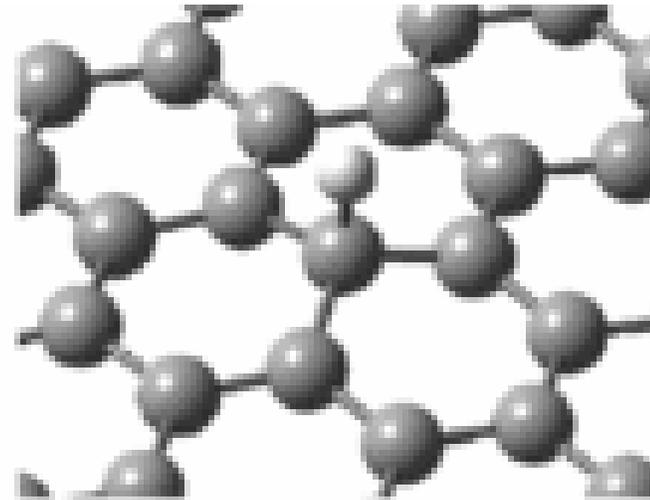
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# DEFECTS INDUCED MAGNETISM

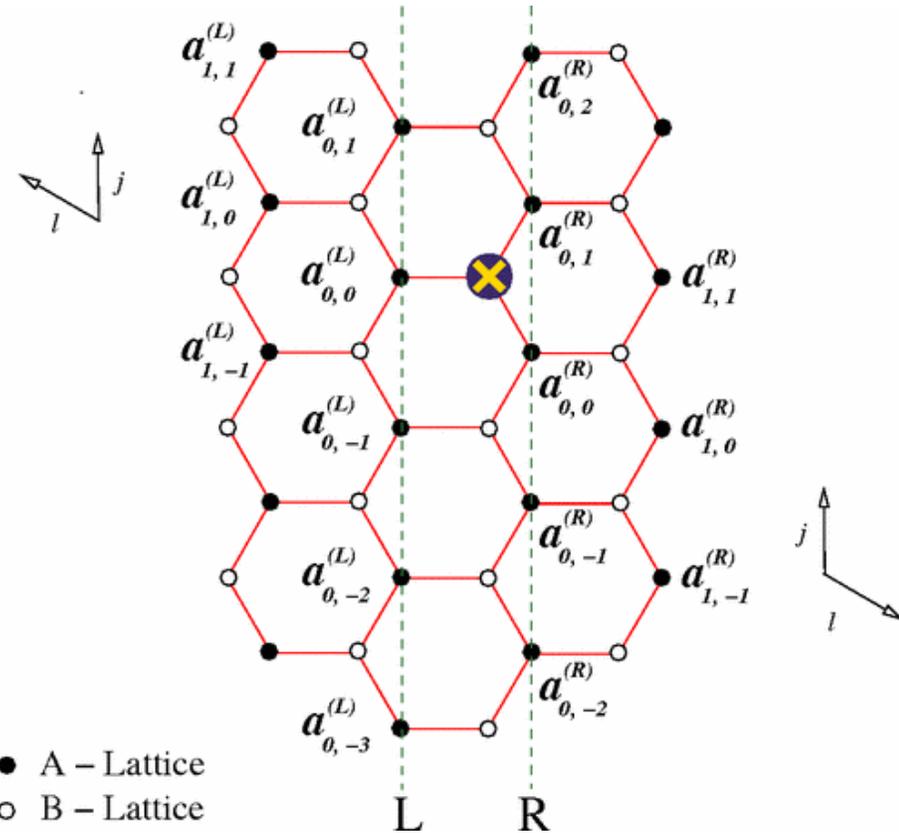
- Imbalance in sublattices leads to magnetism  $N_I \equiv |N_A - N_B| \geq 0$
- Simple defects in graphene

- Hydrogenated carbon:  $sp^2 \rightarrow sp^3$
- Removal of the  $p_z$  orbital
- Protrusion of the carbon
- Acts like a vacancy

HYDROGEN CHEMISORPTION



# QUASI-LOCALISED STATES NEAR VACANCIES

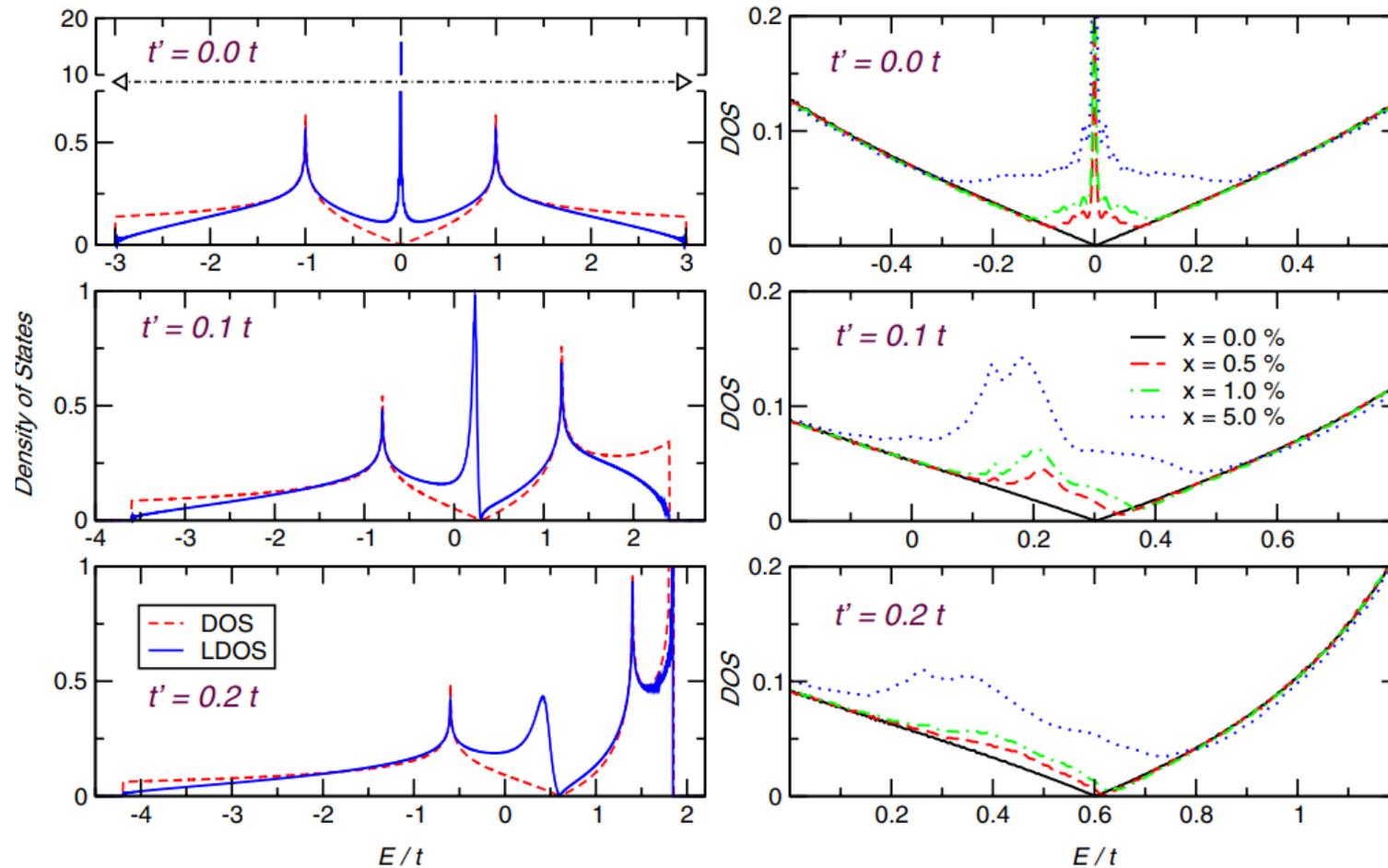


$$\mathcal{H} = -t \sum_{\langle i,j \rangle} c_i^\dagger c_j - t' \sum_{\langle\langle i,j \rangle\rangle} c_i^\dagger c_j + h.c.$$

- $t'$  is the next-nearest-neighbour hopping energy
- ZES are found
- $\Psi_{ZES}(r)$  is quasi-localised
- Power law decay:  $\Psi_{ZES}(r) \propto 1/r$

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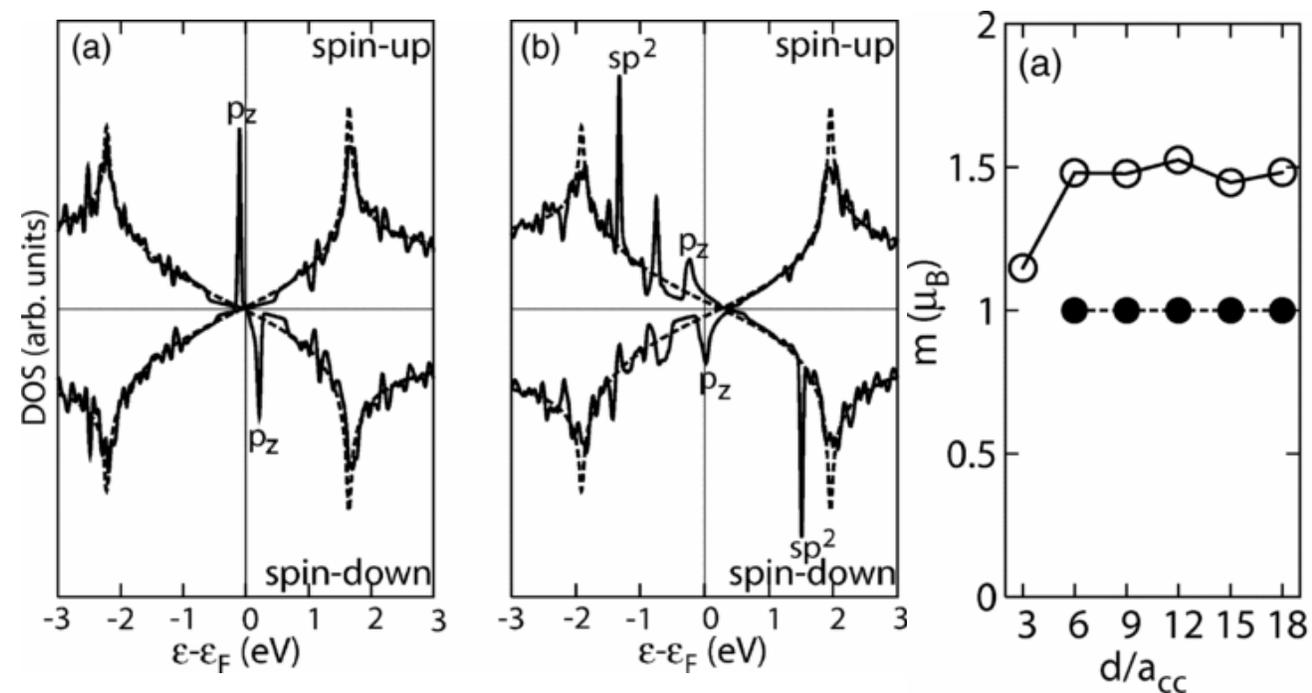
# LDOS OF QUASI-LOCALISED STATES



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# ONSET OF MAGNETISM

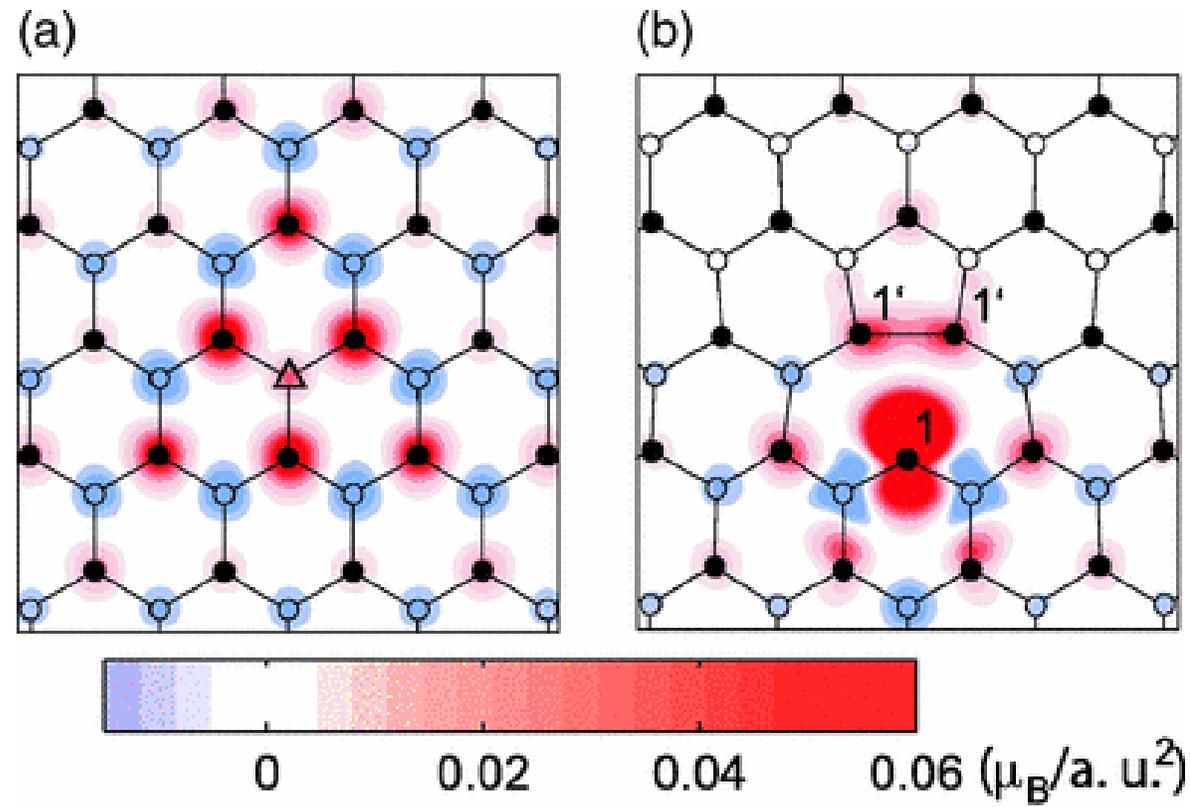
- DFT calculations on a periodical 2D superlattice of defects
- Magnetism from both chemisorption and vacancies



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# ONSET OF MAGNETISM

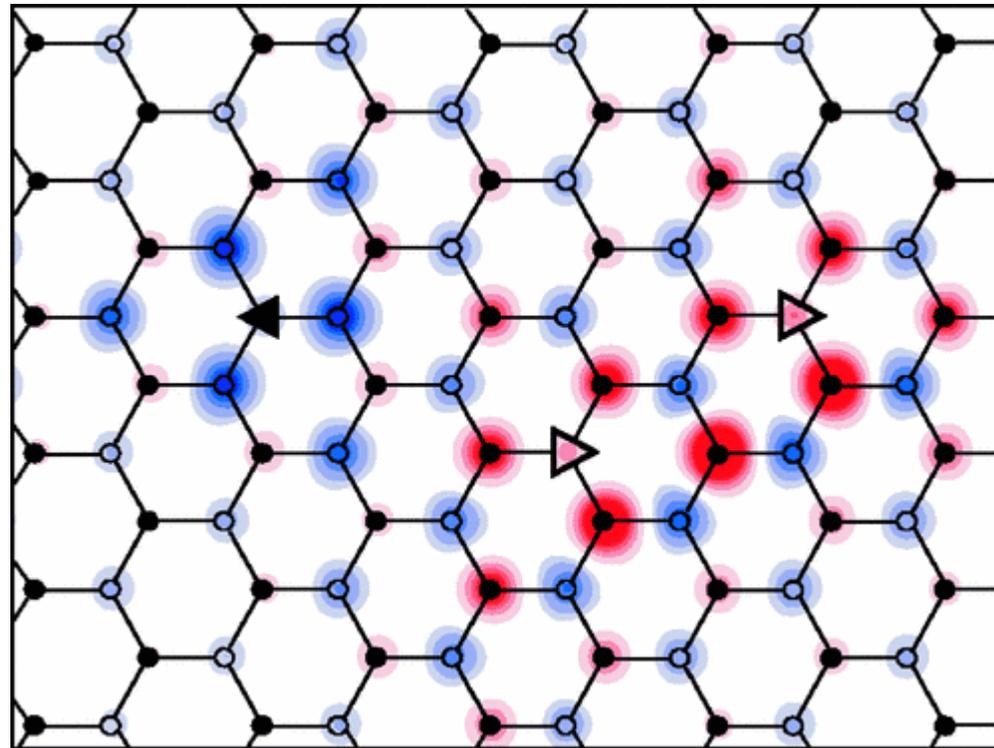
- Spin distribution near the defects



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# ONSET OF MAGNETISM

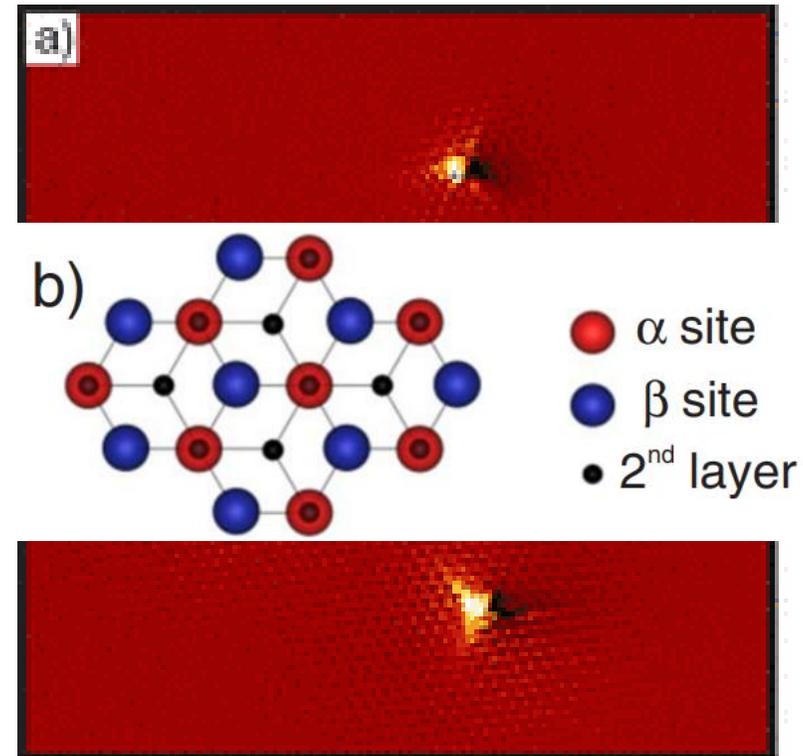
- Studying defects on both sublattices



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# EXPERIMENTAL EVIDENCE

- HOPG with AB Bernal stacking
- Vacancies in exfoliated HOPG by irradiating with 140 eV  $Ar^+$  ions
- Annealing the sample gave single vacancies
- LT-STM was used to study the LDOS



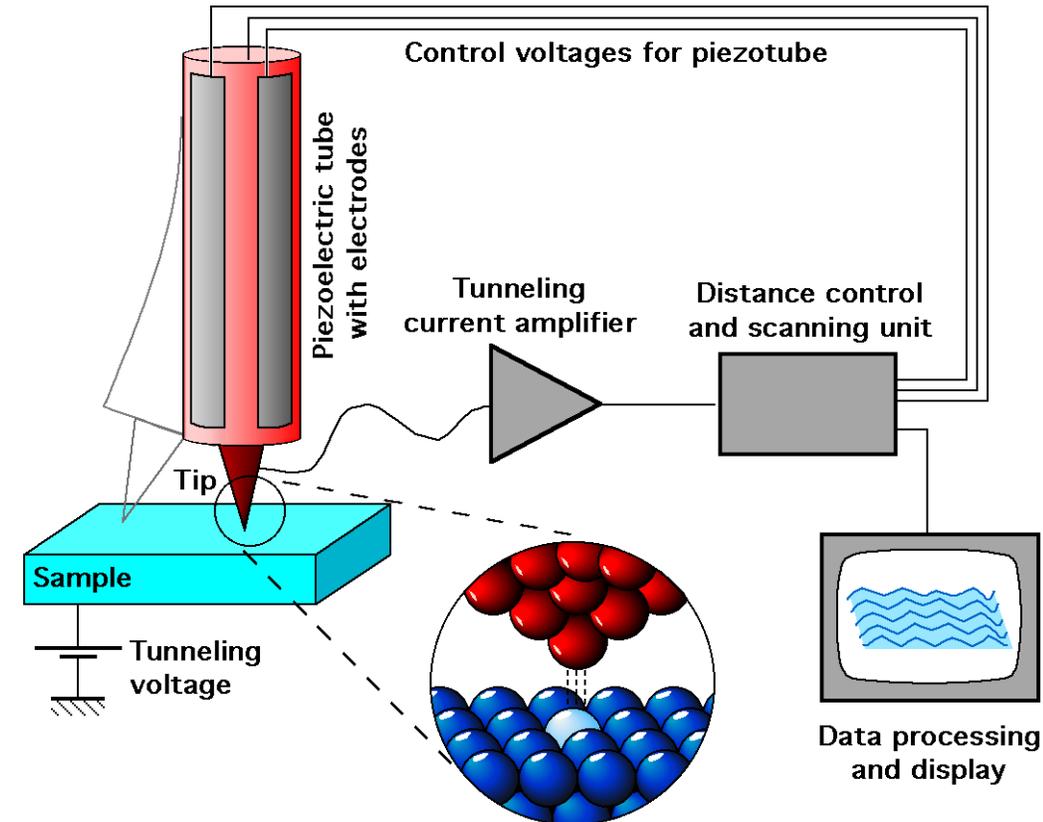
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# LT-SCANNING TUNNELING MICROSCOPY

- STM is based on quantum tunneling
- Measurements done in UHV and LT
- Measures the LDOS

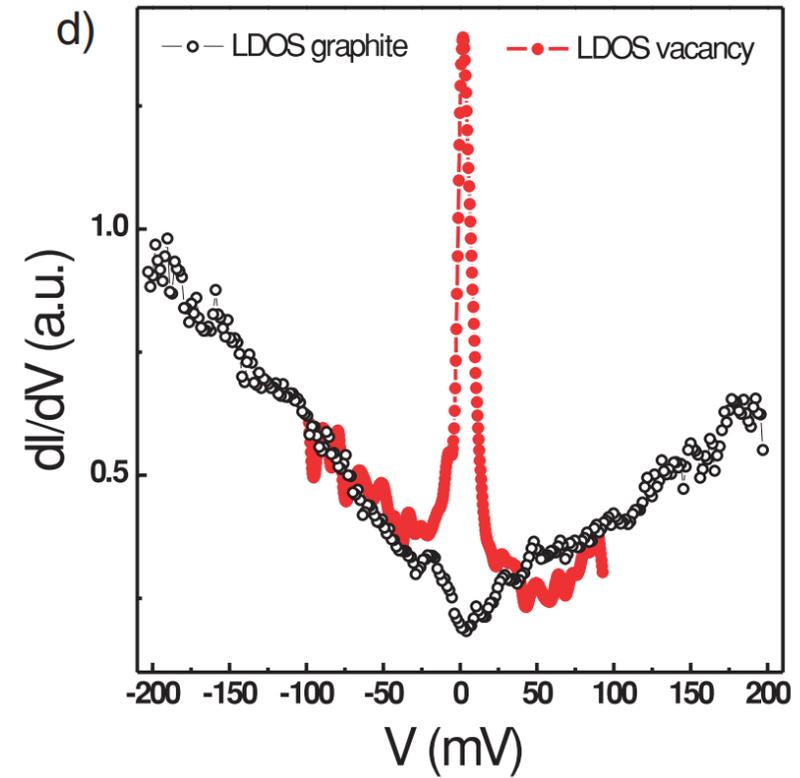
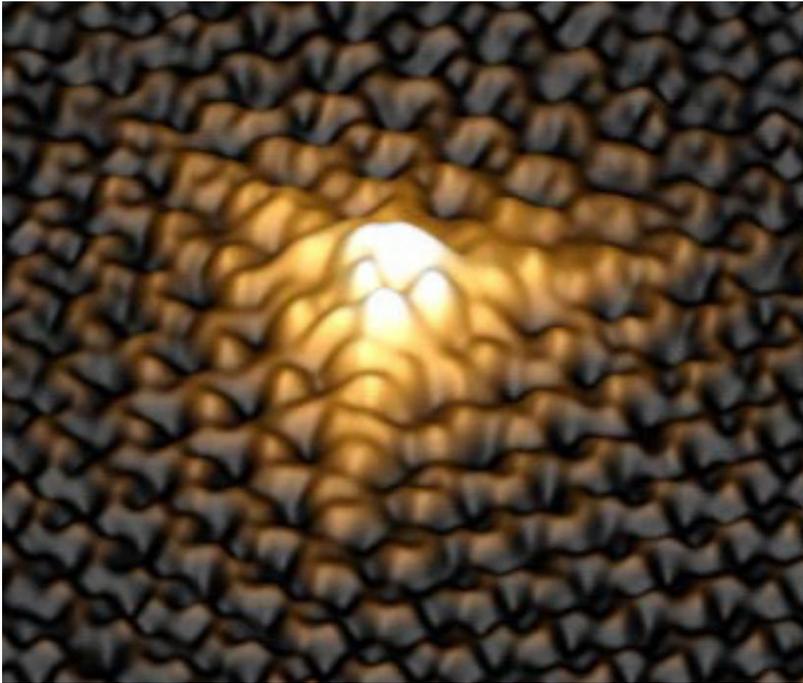
$$I_t = \int_0^{eV} \rho_{\text{sample}}(E, r_0) \rho_{\text{tip}} |M|^2 dE$$

$$\frac{dI_t(V)}{dV} \propto \rho_{\text{sample}}(E = eV, r_0)$$



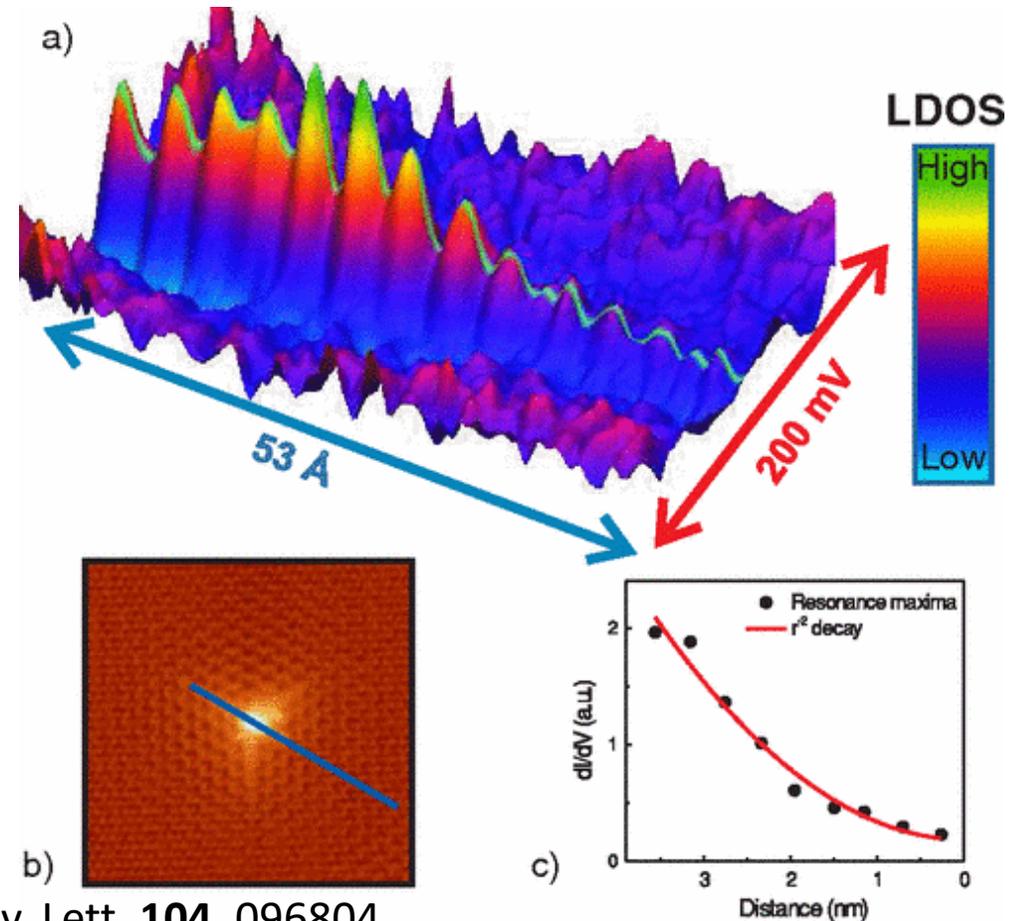
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# LT-SCANNING TUNNELING MICROSCOPY



# EXPERIMENTAL EVIDENCE

- Local LDOS map around a single vacancy
- Shows a peak around  $E_F$
- Peaks maxima vs position is consistent with  $r^{-1}$  decay

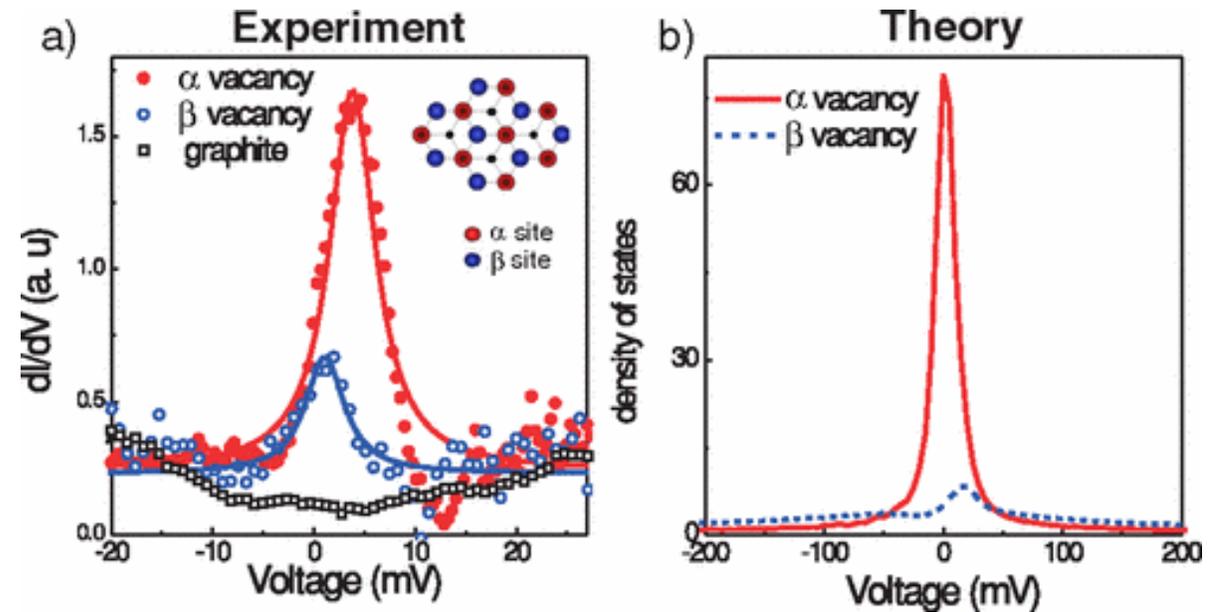


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# AB BERNAL STACKING

- Sublattices are not equivalent
- LDOS peak changes
- In agreement with theory prediction
- Ferrimagnetism arise
- Curie Temperature can be estimate

$$T_C \approx e^2 \sqrt{n_v}$$

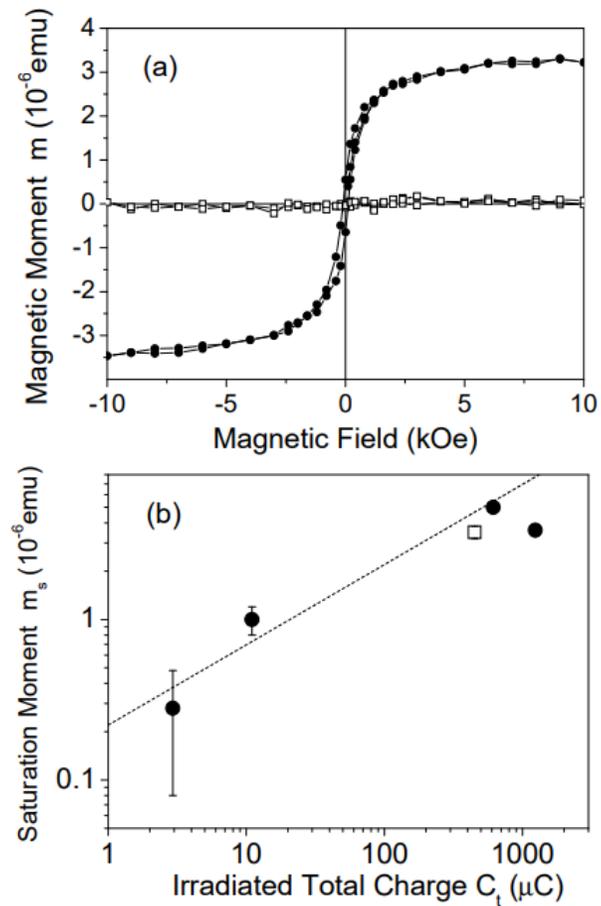
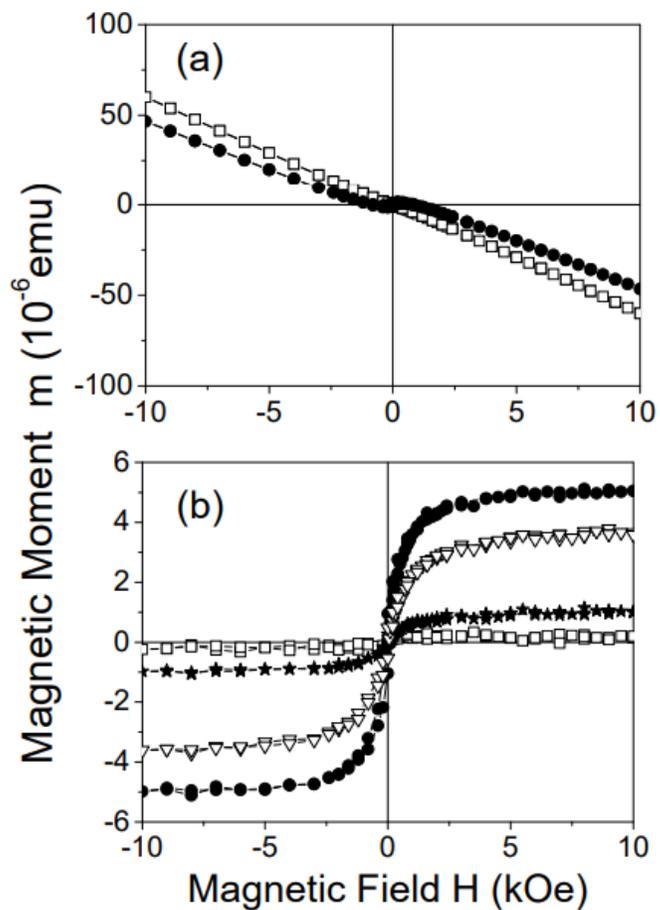


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# BULK DEFECTS

- HOPG irradiated with 2.25 MeV protons
  - Different types of defects
  - Concentration of defects changes with number of irradiation cycles
  - Hysteresis loops were measured at different temperatures 5 K, 300 K, 380 K
  - Ferromagnetic at RT
  - Curie Temperature above 400K
-

# BULK DEFECTS

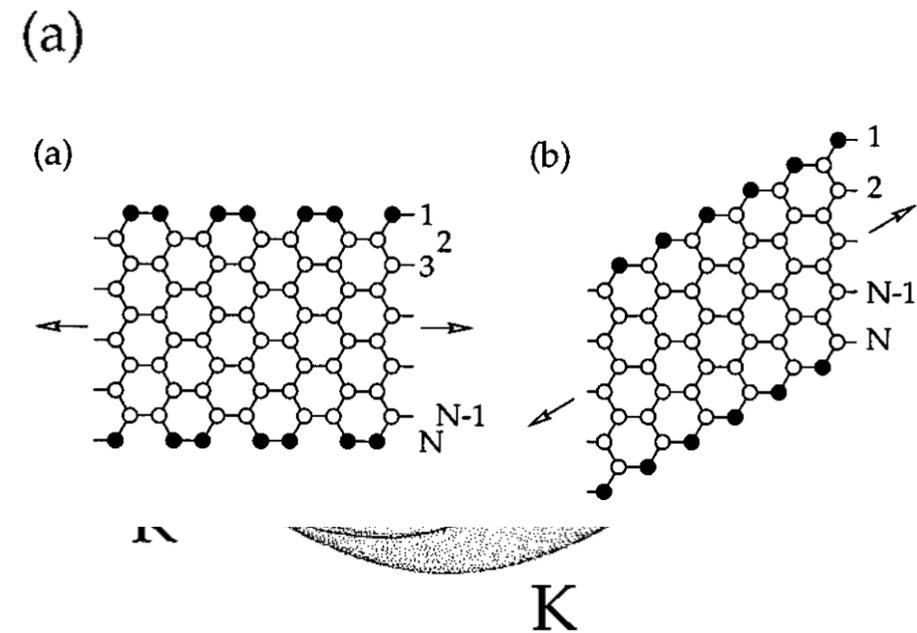


Esquinazi P. et al. 2003 Phys. Rev. Lett. 91 227201

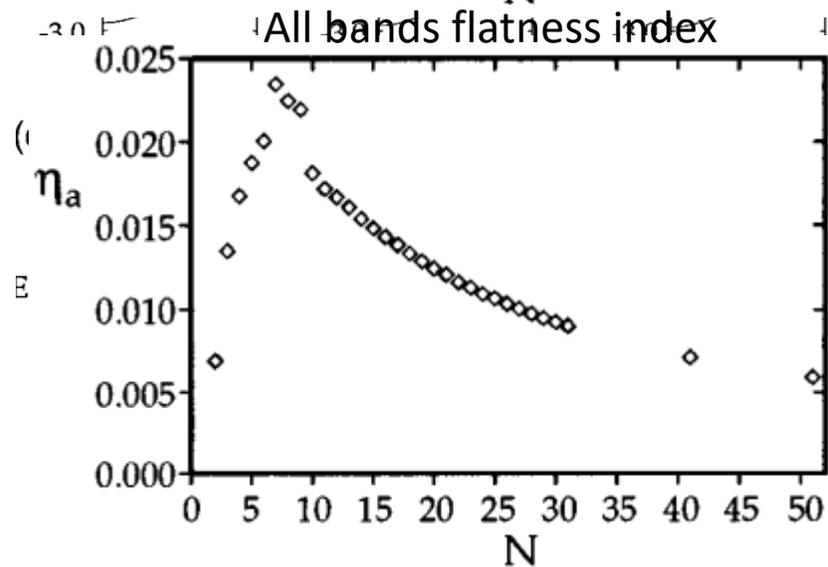
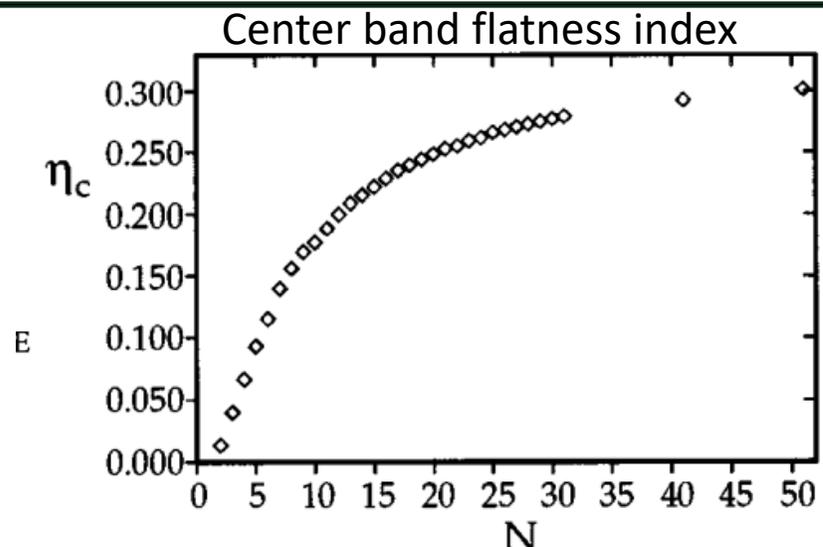
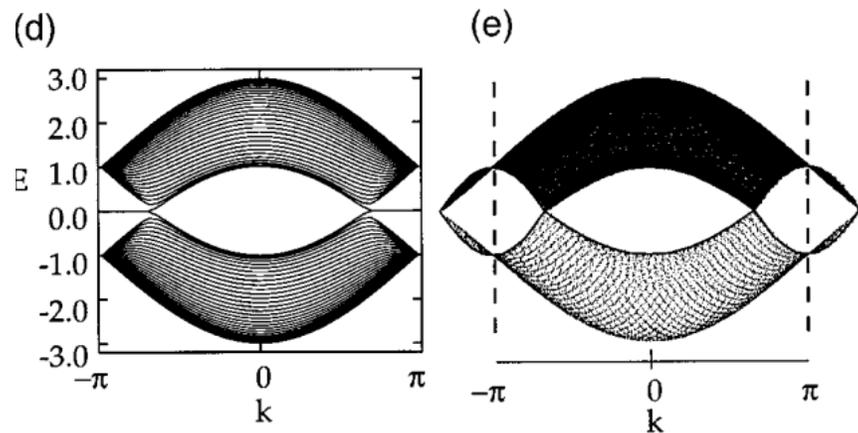
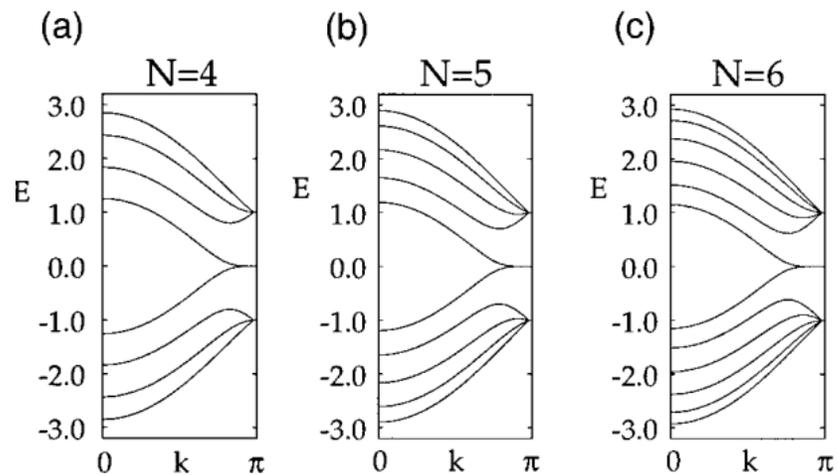
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# EDGE MAGNETISM

- Graphene is a zero-gap semiconductor
- Nanometric size changes the band structure
- Graphene nanoribbons: finite-size 1D system
- Two possible edges
- Magnetism decreases with the increase of size



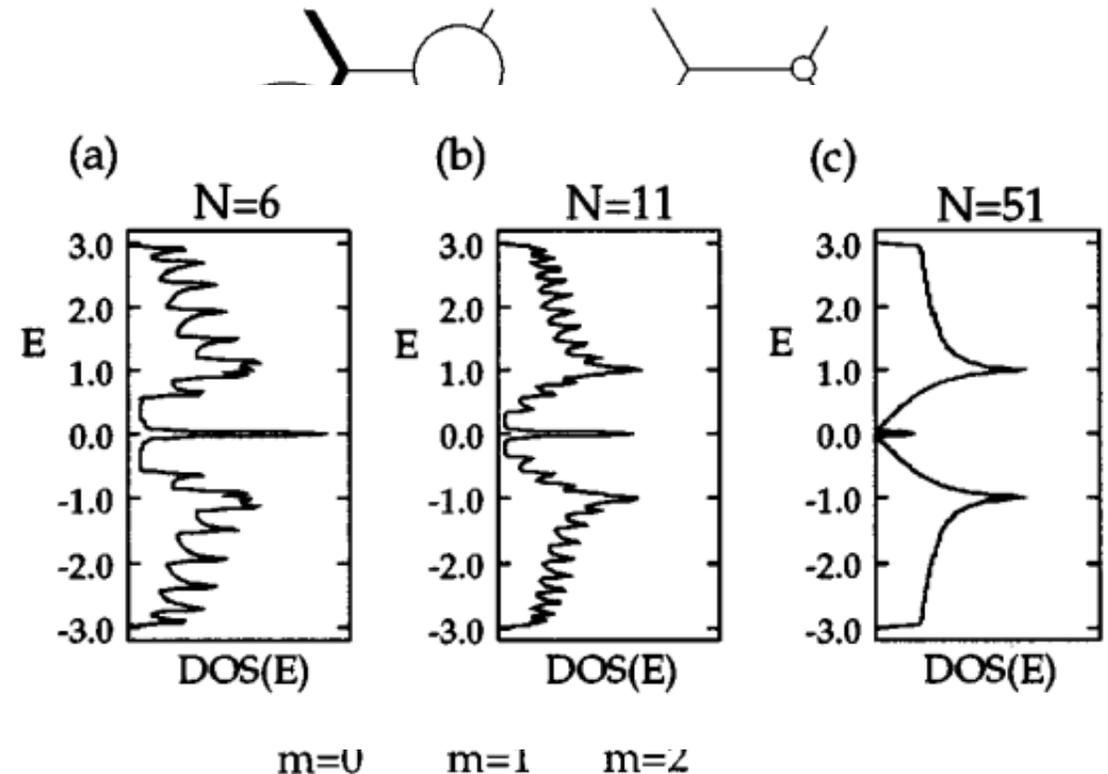
# EDGE MAGNETISM



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# EDGE MAGNETISM: THEORY

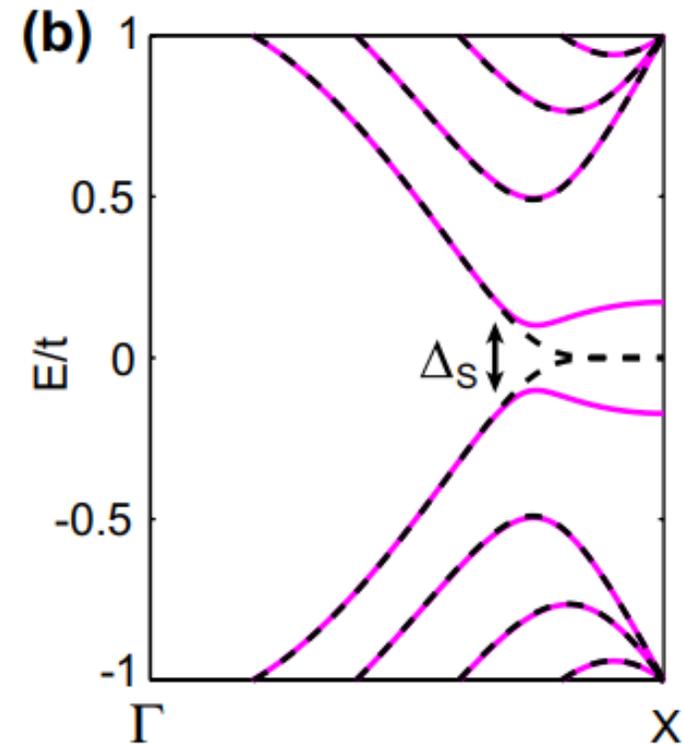
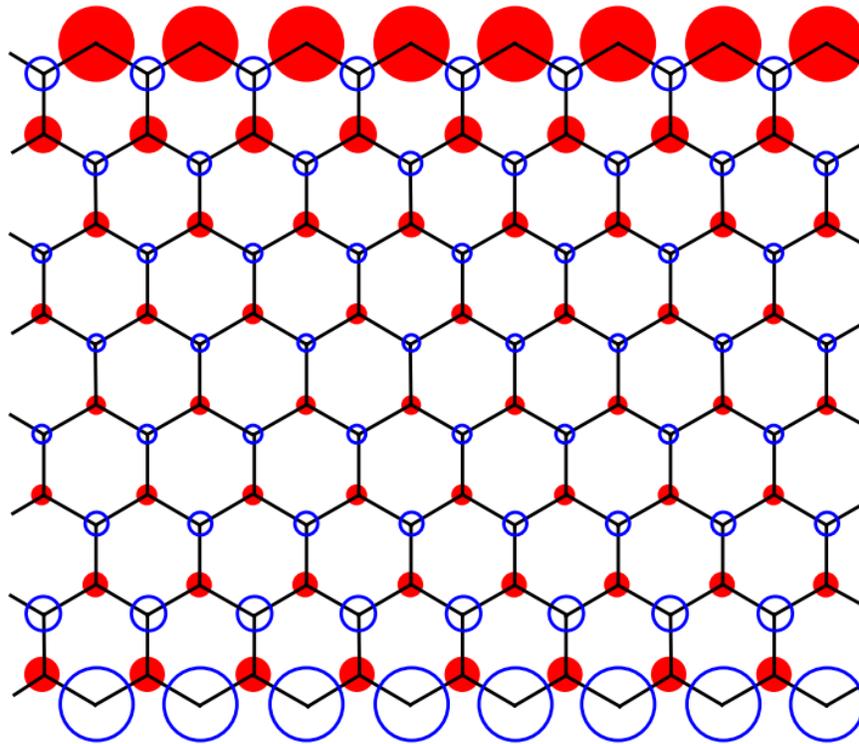
- Semi-infinite graphene plane with zig-zag edge
- Tight-binding of LCAO is used
- Spatially localized states are found
- States localised on a single sublattice
- $\psi_s(r)$  has an exponential decay



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# EDGE MAGNETISM: THEORY

Considering the interaction between electrons using the Hubbard model



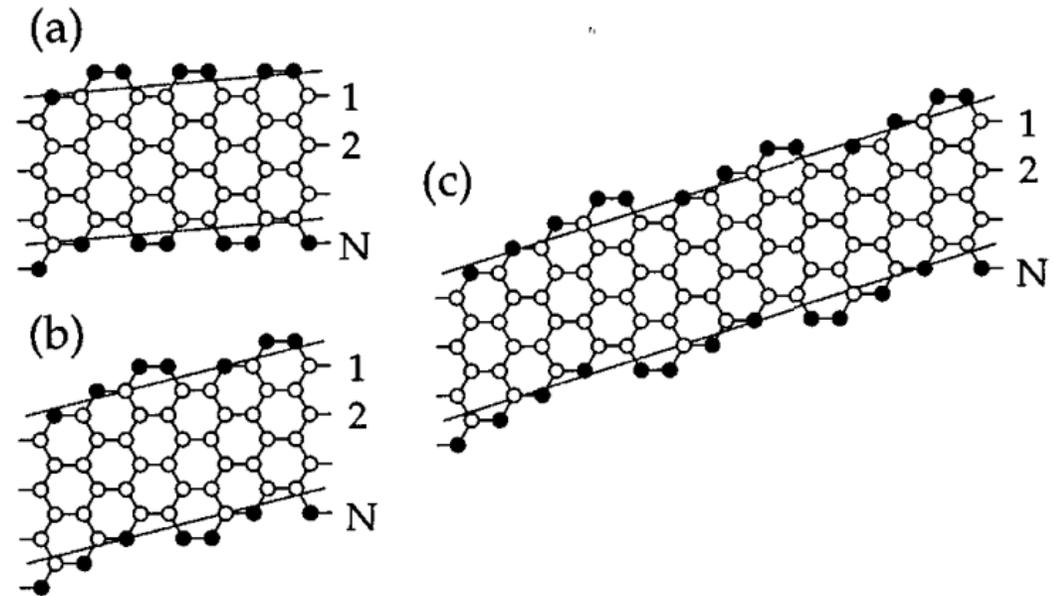
Yazyev O V 2008b Phys. Rev. Lett. 101 037203

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# EDGE MAGNETISM: REAL EDGES

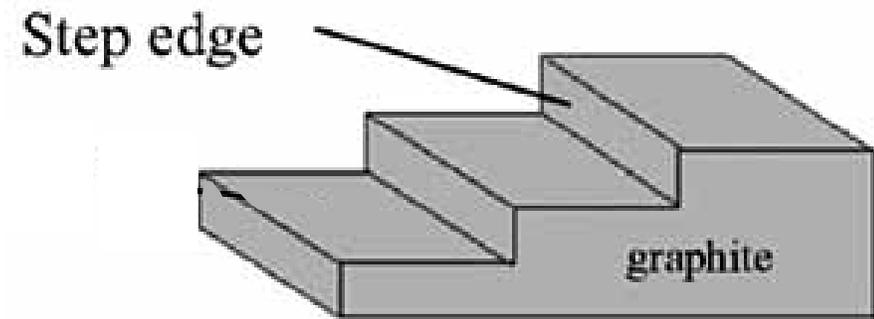
- Real edges are of both types
- The ratio can be different
- Spatially localized states are still found



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# EDGE MAGNETISM: EXPERIMENTS

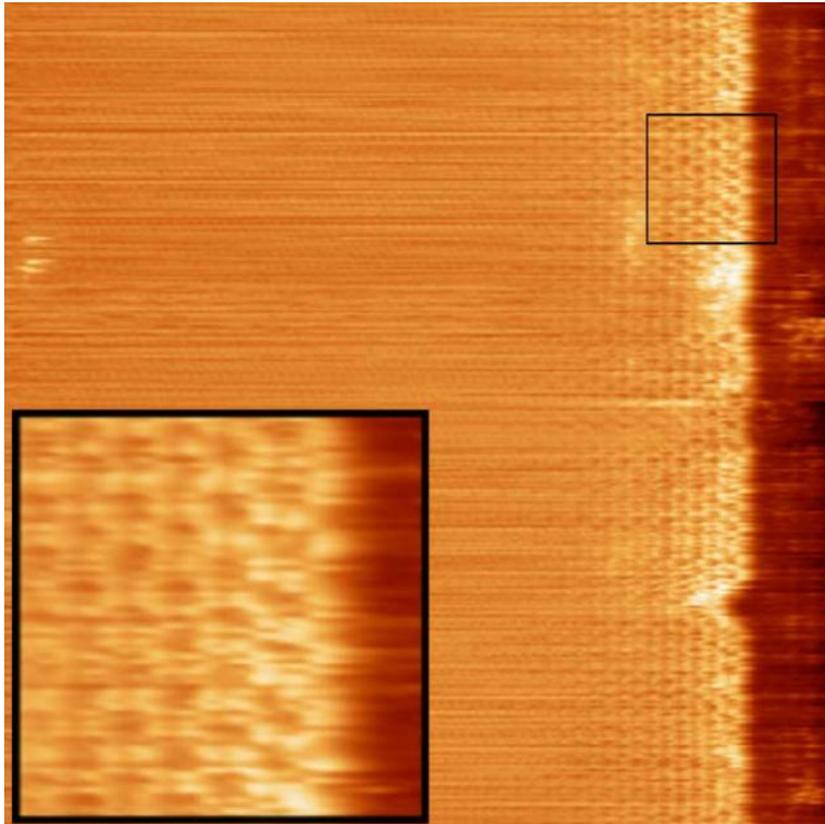
- UHV-STM was used on step edges of HOPG
- Brightness of the images is proportional to LDOS
- Theoretical predictions were confirmed



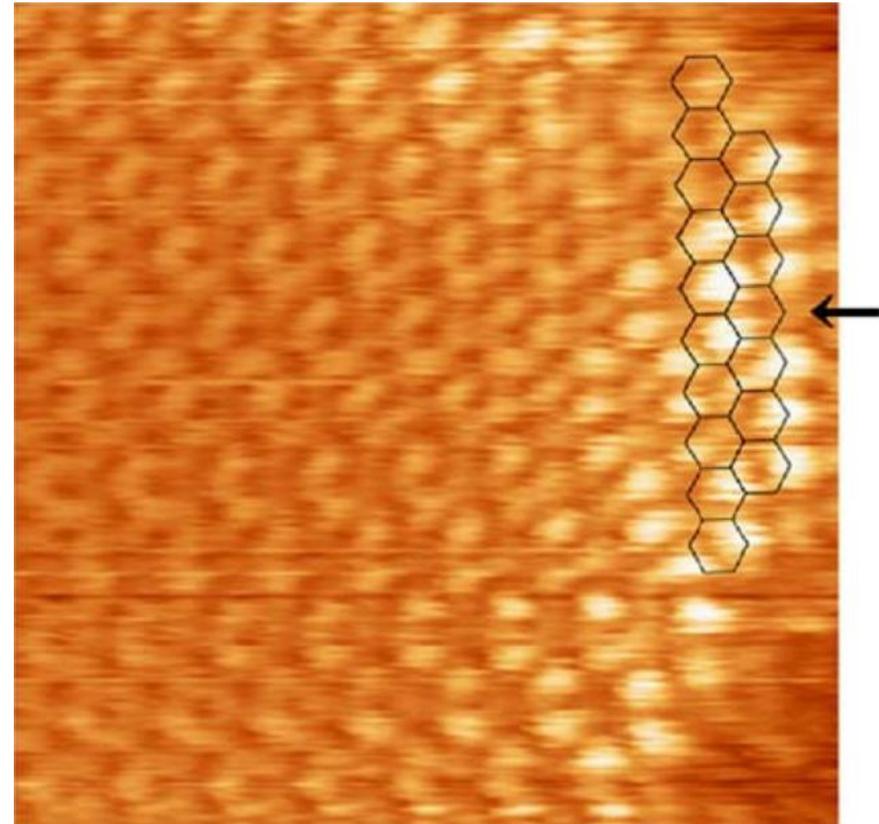
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# EDGE MAGNETISM: EXPERIMENTS

Armchair edge



Zigzag edge

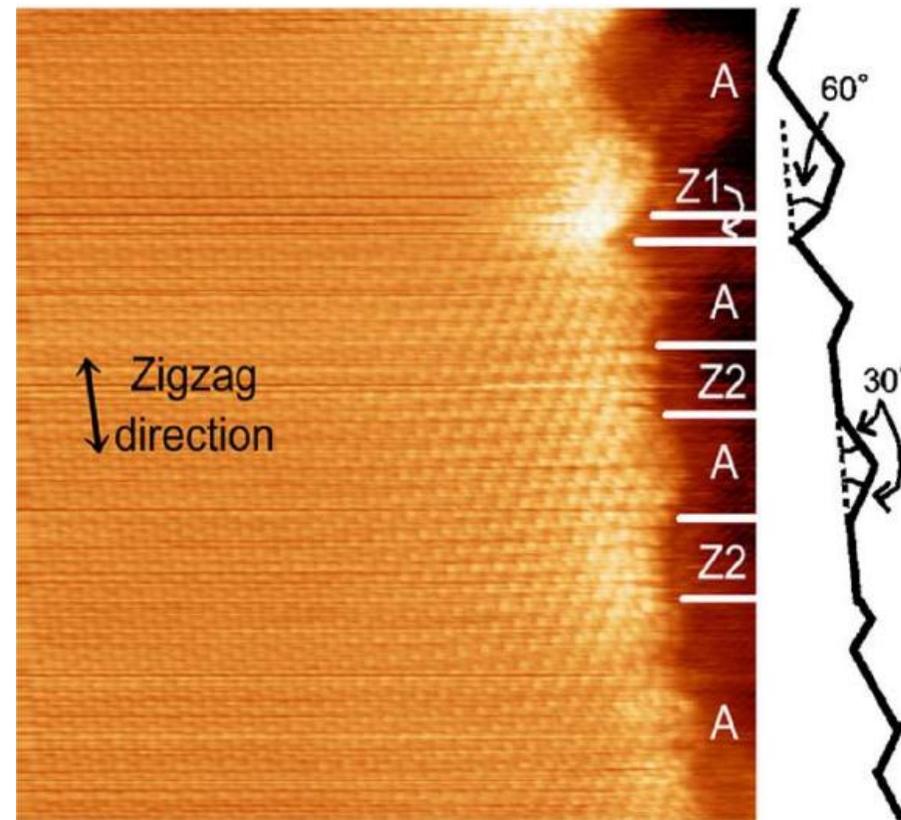
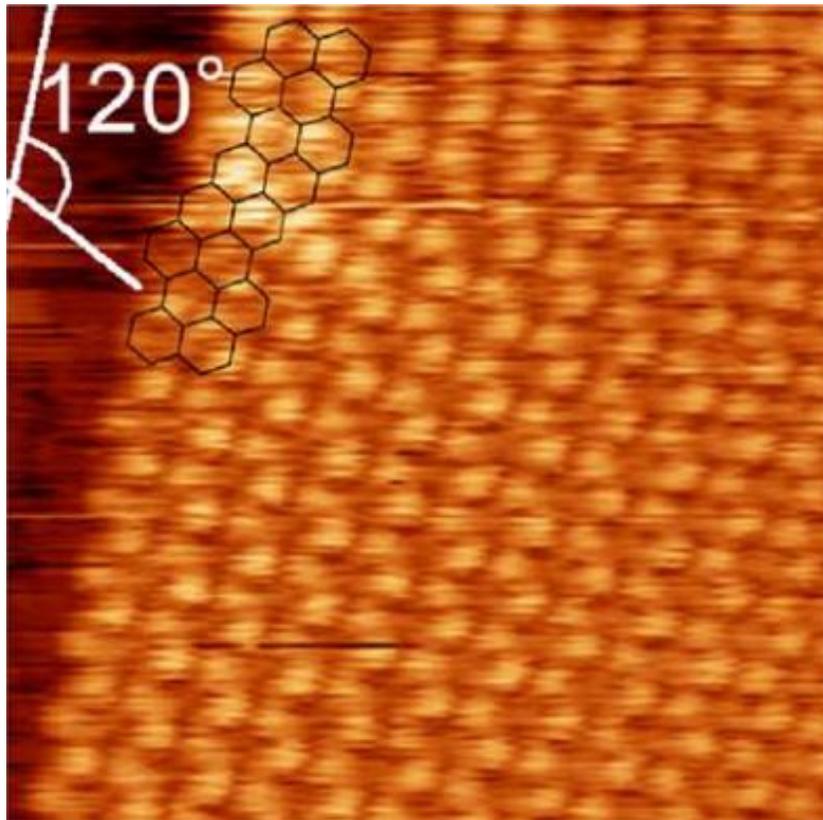


Kobayashi Y, Fukui K i, Enoki T and Kusakabe K 2006 Phys. Rev. B 73 125415

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# EDGE MAGNETISM: EXPERIMENTS



Kobayashi Y, Fukui K i, Enoki T and Kusakabe K 2006 Phys. Rev. B 73 125415

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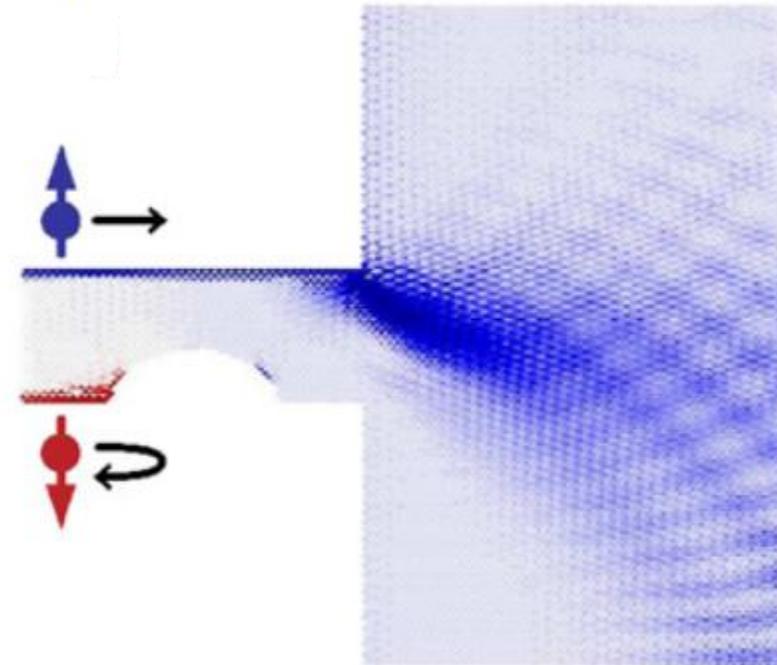
# **APPLICATIONS**

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# APPLICATIONS: SPINTRONIC

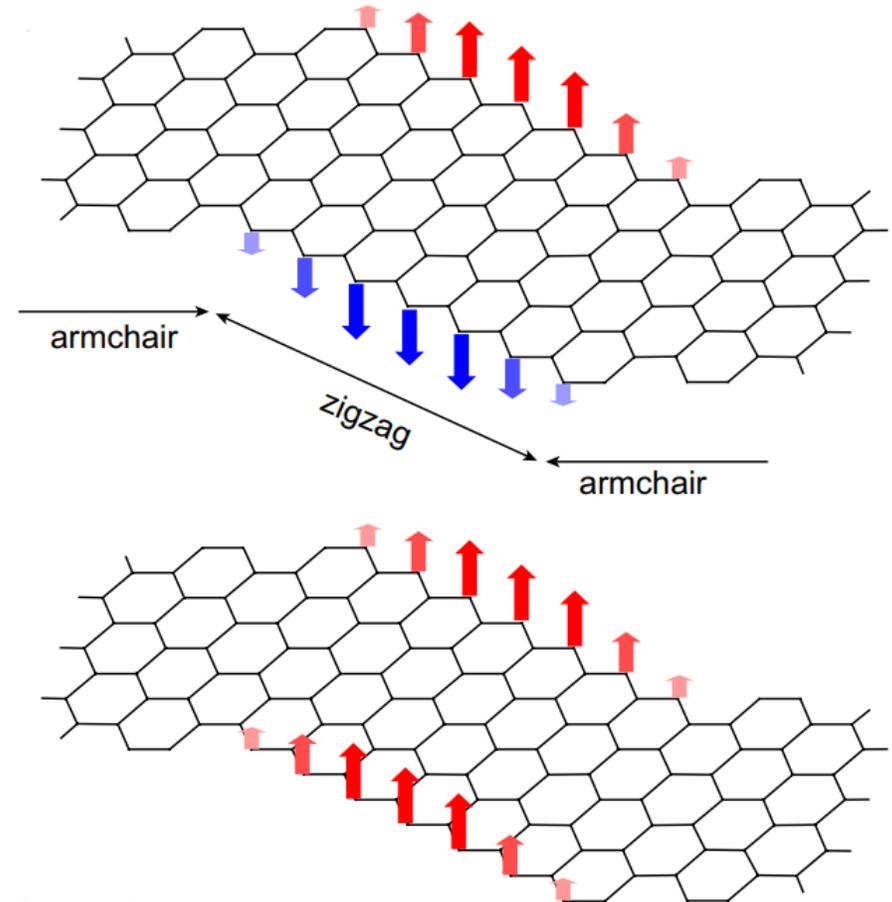
- Electrically or defect induced half-metallicity in zigzag graphene nanoribbons
- Nanoribbon attached to a graphene layer
- Spin-polarised electrons are injected into the graphene sheet
- An extended defect is introduced



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# APPLICATIONS: MAGNETORESISTIVITY

- Graphene nanostructure in magnetoresistive junctions
- High MR
- An all graphene-based device
- Low resistivity state with magnetic field
- High resistivity without



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## REFERENCES

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  - [2] Yazyev O V 2008b Phys. Rev. Lett. 101 037203
  - [3] Kobayashi Y, Fukui K i, Enoki T and Kusakabe K 2006 Phys. Rev. B 73 125415
  - [4] Esquinazi P, Spemann D, H"ohne R, Setzer A, Han K H and Butz T 2003 Phys. Rev. Lett. 91 227201
  - [5] V. M. Pereira, F. Guinea, J. M. B. Lopes dos Santos, N. M. R. Peres, and A. H. Castro Neto, Phys. Rev. Lett. 96, 036801 (2006).
  - [6] M. M. Ugeda Phys. Rev. Lett. **104**, 096804
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